

Simplest Neutral Singlet C₂E₄ (E = Al, Ga, In, and Tl) Global Minima with Double Planar Tetracoordinate Carbons: Equivalence of C₂ Moieties in C₂E₄ to Carbon Centers in CAI₄²⁻ and CAI₅⁺

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Received: November 11, 2008; Revised Manuscript Received: January 13, 2009

Ab initio and DFT calculations have been carried out to search for the simplest neutral singlet species with double planar tetracoordinate carbons (dptCs) [the “simplest” means the species containing the least number (six) and types (two) of atoms]. Under the restrictions to the possible models (M1–M4) with dptCs and to the singlet electronic states, the B3LYP/6–31+G* scanning on the candidates, C₂E₄ (E = the second- and third-row main group elements), only led to two minima (*D*_{2h} C₂Al₄ and *C*_{2h} C₂Be₄) with stable DFT wave functions. The extensions to the heavier elements after the fourth row in the IIA and IIIA groups revealed that the *D*_{2h} C₂E₄ (E = Ga, In, and Tl) are also minima with dptCs but C₂Ca₄ (*C*_{2h}) is a first-order saddle point. Extensive explorations at the DFT level on their potential energy surfaces (PESs) further confirmed that the *D*_{2h} C₂E₄ (E = Al, Ga, In, and Tl) are the global minima, but the *C*_{2h} C₂Be₄ is a local minimum. The optimizations at the MP2 level distorted the *D*_{2h} C₂E₄ (E = Ga, In, and Tl) slightly and the distortion energies are less than 0.02 kcal/mol. The C₂E₄ (E = Al, Ga, In, and Tl) with dptCs are 18.0, 18.3, 13.4, and 12.2 kcal/mol energetically more favorable than their nearest isomers, respectively, at the CCSD(T)//MP2 level with aug-cc-pVTZ for C and Al and aug-cc-pVTZ-PP for Ga, In, and Tl basis set. The substantial energy differences suggest their promise to be experimentally realized. The strong peak on the C₂Al₄⁻ component in the time-of-flight mass spectrum from laser vaporization of a mixed graphite/aluminum may relate to the *D*_{2h} C₂Al₄ global minimum. The analyses of the electronic structures of C₂Al₄ (*D*_{2h}), CAI₄²⁻ (*D*_{4h}) and CAI₅⁺ (*D*_{5h}) indicates that the C₂ moiety in C₂Al₄ is the equivalence of carbon centers in CAI₄²⁻ and CAI₅⁺ and unveils the reasons for their stability. The electronic structures of C₂Al₄ and ethene are compared. On the one hand, an Al atom functions like an H atom because the eight more valence electrons of C₂Al₄ than C₂H₄ occupy four nonbonding orbitals and are not effectively utilized for bonding. On the other hand, an Al atom is different from an H atom because an Al atom has p electrons available for peripheral bonding around the C₂ moieties in C₂Al₄, which further rationalize the origins for C₂E₄ to achieve double ptCs.

1. Introduction

Linear dicoordination (e.g., ethyne), planar trigonal tricoordination (e.g., ethene) and tetrahedral tetracoordination (e.g., methane) are the predominant bonding patterns of carbon. However, a curious exception, planar tetracoordinate carbon (ptC), has fascinated chemists since the 1970s, when Hoffmann, Alder, and Wilcox¹ suggested ways to reduce the unfavorable energies of ptC disposition after Monkhorst proposed the planar configuration for the interconversion of enantiomers with asymmetric carbon in 1968.² Since the first computational prediction of a ptC minimum, 1,1-dilithiocyclopropane, by the Schleyer–Pople group,³ various molecules with ptCs have been characterized computationally or experimentally.^{4–9} Because the coordination number for the conventional planar carbon bonding (e.g., the carbons in ethene and benzene) is three, one can consider planar carbon having a coordination number higher than three, as planar hypercoordinate carbon (phC). Molecules with planar penta- and hexa-coordinate carbon have been computationally predicted to be viable by Wang and Schleyer¹⁰ and Exner and Schleyer,¹¹ respectively. Recently, the planar

bonding patterns have been further extended to planar hypercoordinate heteroatoms and transition metals.^{12–18}

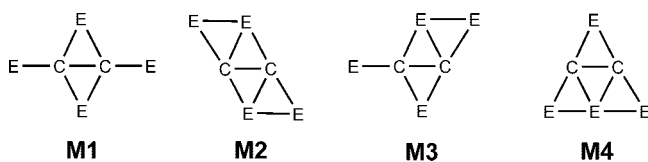
In the phC chemistry, chemists, on one hand, strive to design species containing more and more phCs. For example, the Frenking–Schleyer group¹⁹ recently reported boron rings with multiple phCs. Minyaev et al.^{20,21} designed species with several ptCs based on C₃B₂H_n (*n* = 2, 4, and 6) ptC units. Zhang and Cao^{22,23} constructed zigzag C–B nanotube with quasi-ptCs. Using the CM₄H₄ (M = Ni, Pd, Pt) as basic building blocks, Wu et al.²⁴ designed compounds of various shapes with multiple ptCs or quasi-ptCs. On the other hand, they make efforts to achieve phC in the simplest global minima because being the global minimum can facilitate experimental generations and identifications. For example, the *D*_{4h} CAI₄²⁻ has been identified to be the simplest penta-atomic global minimum, this species and its analogs, CAI₄⁻,²⁵ CAI₃Si⁻,²⁶ CAI₃Ge⁻,²⁶ CAI₄²⁻, and NaCAI₄⁻,²⁷ have been experimentally detected in 1999 and 2000. Interestingly, their isoelectronic ptC species, *cis*- and *trans*-CAI₂Si₂, were predicted by Schleyer and Boldyrev in 1991. However, the effort to detect the phC species CB₇⁻ and CB₆²⁻ (not global minima)^{10,11} led to the more energetically favorable species, in which the carbon atoms prefer locating on the edges or vertexes.^{28,29} More examples of simplest global minima with phC, such as CB₄,^{30,31} CAI₅⁺,³² CCu₄²⁺, and its isoelectronic analogs,³³ have recently been reported. Moreover, the simplest species may be used as the basic building blocks to construct new molecules or materials with multiple phCs. The potential

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SCHEME 1: Simplest Models Possibly Containing dptCs



of using small ptC building blocks (e.g., CB_6^{2-} ,^{34–39} C_5^{2-} ,^{40–45} CAI_3Si ,⁴⁶ and CAI_4^{2-} ,⁴⁷) to design novel materials has been explored theoretically. In this work, we attempted to find the simplest species with double planar tetracoordinate carbons (dptCs), and showed that the obtained C_2E_4 (E = Al, Ga, In, and Tl) are not only the simplest species with dptCs, but also the global minima. Note that the “simplest” means the dptC species containing the least number (six) and types (two) of atoms.

2. Computational Methods

Scheme 1 sketches the simplest binary models (M1–M4) which possibly contain dptCs. Under the restrictions to these models and to singlet electronic states, we scanned the second- and third-row main group elements at the B3LYP/6–31+G* level. The species, C_2Al_4 (M1) and C_2Be_4 (M2), are the only survivors to meet two criteria: minima with dptCs and the DFT wave functions being stable. Other species, such as C_2Li_4 (M1 and M2) and C_2Na_4 (M2 and M4), are minima with dptCs but their wave functions are unstable. C_2Mg_4 has no dptCs minimum under the restrictions. We then extended the scanning to the heavier elements in the group of IIA and IIIA at the B3LYP level with basis sets 6–31+G* for C, Ca, and Ga and aug-cc-pVDZ-PP^{48–50} for In and Tl. At this level, C_2E_4 (E = Ga, In, and Tl) are also minima but C_2Ca_4 is a first-order saddle point. We therefore discarded C_2E_4 (E = Mg and Ca) in the following.

The located species with dptCs were then verified to be global minima using the following procedure. For each C_2E_4 (E = Be, Al, Ga, In, and Tl), 3000 random structures were generated by a program called GXYZ,⁵¹ which is principally the same as the “kick” method proposed by Saunders et al.^{52,53} Both of the methods generate a large pool of random structures for geometric optimizations. About 80~90% of these random structures were filtered out because they are structurally unreasonable. The remaining 376, 278, 202, 171, and 185 structures for E = Be, Al, Ga, In, and Tl, respectively, were then subjected to optimizations at B3LYP level with the basis

set mentioned above. Ignoring the slight differences due to releasing the symmetry constraint, there were 6, 8, 11, 10, and 12 structures for E = Be, Al, Ga, In, and Tl, respectively, identified to be the same as the dptCs minima optimized on the basis of the models. This implies the good samplings on the potential energy surfaces (PESs). To further examine the sampling convergence, taking C_2Al_4 as an example, another independent exploration with 4000 initial random structures was carried out. Optimizations on the 717 initially filtered input structures gave 40 structures having D_{2h} symmetry with dptCs. In a study on the C_nAl_m ($n = 2–3$, $m = 2–8$) clusters, Naumkin⁵⁴ also reported the D_{2h} C_2Al_4 species, but it is unclear if it is the global minimum. The dptCs minima, together with the lowest three for C_2E_4 (E = Al, Ga, In, and Tl) or five for C_2Be_4 isomers obtained in the above step, were then refined and confirmed to be minima by frequency analysis calculations at the B3LYP/aug-cc-pVTZ (aug-cc-pVTZ-PP^{48–50} for Ga, In, and Tl) level. Concerning the caveats about the DFT reliability recently raised by several authors,^{55–59} the B3LYP geometries (including the dptCs minima) were reoptimized at the MP2/aug-cc-pVTZ (aug-cc-pVTZ-PP for Ga, In, and Tl) level, which resulted in slightly distorted structures for E = Ga, In, Tl (see below). The energetic results were finally improved at the CCSD(T)/aug-cc-pVTZ (aug-cc-pVTZ-PP for Ga, In, and Tl) level using the MP2 structures. The results reported in the following are all calculated with the basis set aug-cc-pVTZ (aug-cc-pVTZ-PP for Ga, In and Tl). The basis set will not be mentioned hereafter, unless otherwise specified. All ab initio and DFT calculations were performed using the *Gaussian 03* program of package.⁶⁰ The electronic structures were analyzed by NBO 5.0^{61,62} at the B3LYP level.

3. Results and Discussion

Being the structures of C_2E_4 (E = Al(1a), Ga(2a), In(3a), and Tl(4a)), optimized at the B3LYP and MP2 levels, are shown in Figure 1, and their Cartesian coordinates at the two levels are given in the Supporting Information I (SI1). At the B3LYP level, the D_{2h} 1a–4a are all minima, but at the MP2 level, only D_{2h} 1a is a minimum and D_{2h} 2a, 3a, and 4a are first-order saddle points with very small imaginary frequencies (9i, 16i, and 10i cm^{-1} , respectively). Following the vibrational modes corresponding to the imaginary frequencies, the D_{2h} 2a, subjecting to geometrical optimization, degenerated to a planar C_{2v} structure with dptCs, but the D_{2h} 3a and 4a to the bending C_{2v} structures without perfect ptC (see Figure 1). However, the

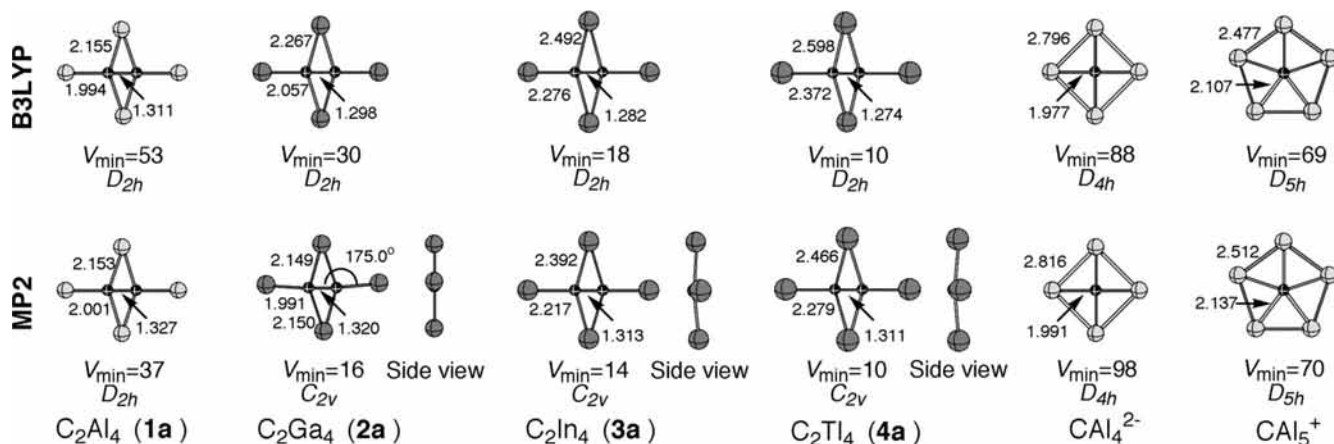


Figure 1. B3LYP and MP2 optimized geometries of 1a–4a, CAI_4^{2-} , and CAI_5^+ , together with the numbers of imaginary frequencies (NIMAG) and the lowest vibrational frequencies (V_{\min}) in cm^{-1} . Bond lengths and bond angles are given in angstroms and degrees, respectively.

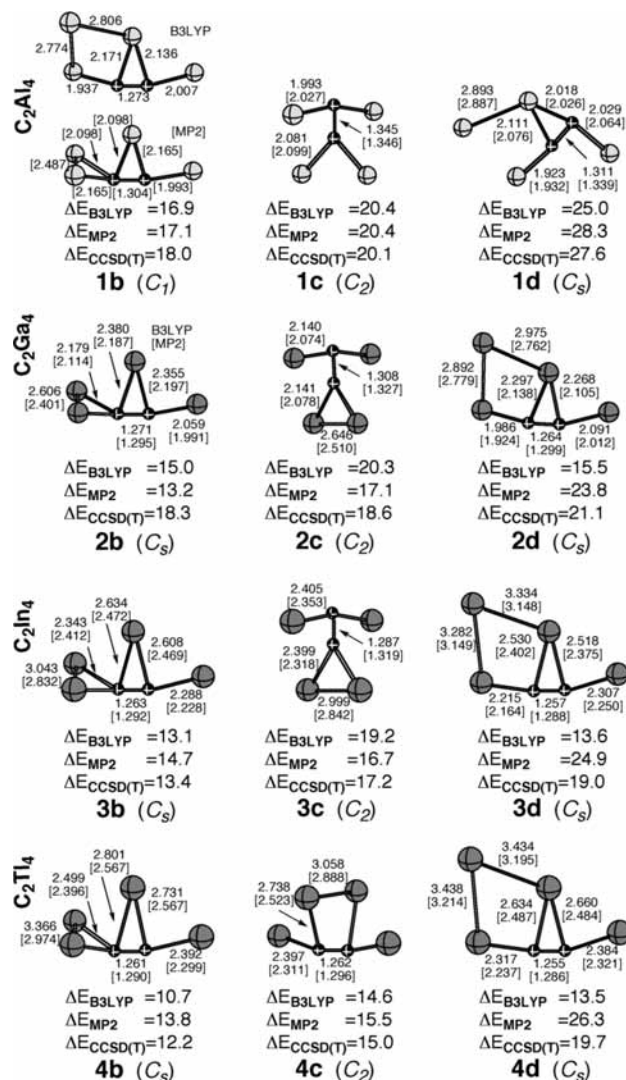


Figure 2. B3LYP and MP2 optimized geometries of the three singlet isomers nearest to **1a–4a**, together with the relative energies (ΔE in kcal/mol) at the B3LYP, MP2, and CCSD(T)//MP2, respectively. Bond lengths are given in Å. Values in brackets are MP2 results.

energy differences between C_{2v} and D_{2h} structures of **2a**, **3a**, and **4a** are very small (less than 0.02 kcal/mol). The negligible differences are much less than the zero point energy corrections and therefore the C_{2v} **3a** and **4a** can be considered to be fluxionary ptC species and detected as vibrationally averaged planar D_{2h} structures experimentally.^{25,63}

The global minimum is very important for feasibly experimental access, in particular for the laser-ablation-based experiments because the method tends to anneal species into the global minimum.⁶³ For example, the experimentally detected D_{4h} CaI_4^- and CaI_4^{2-} have been verified to be the global minima computationally.^{25,27} However, as mentioned before, the effort to detect the phC species CB_7^- and CB_6^{2-} (not global minima)^{10,11} led to the energetically more favorable species, in which the carbon atoms prefer locating on the edges or vertices.^{28,29} The D_{5h} CaI_5^+ ³² is also predicted to be the global minimum. Using the procedure described above, we explored the PESs of C_2E_4 ($E = Al, Ga, In,$ and Tl) extensively. Figure 2 gives the first three lowest isomers of **1a**, **2a**, **3a**, and **4a**. Although with respect to the symmetry planes containing dptCs, the D_{2h} structures are fluxional, they are rigid with respect to other freedoms of motions. At the B3LYP level, **1a–4a** are 16.9, 15.0, 13.1, and 10.7 kcal/mol energetically more favorable

than their nearest isomers (**1b–4b**), respectively. The B3LYP relative energies are reasonably in agreement with the MP2 (17.1, 13.2, 14.7, and 13.8 kcal/mol, respectively) and CCSD(T)//MP2 (18.0, 18.3, 13.4, and 12.2 kcal/mol, respectively) values.

We further considered the triplet states of these species and explored their PESs using the same approaches as for the singlet states. Their structural and energetic results are given in SI1, SI4 and SI5. At our highest CCSD(T)//MP2 level, the lowest triplet states of C_2E_4 ($E = Al, Ga,$ and In) are 1.6, 2.0, and 0.4 kcal/mol higher in energy than their fourth lowest singlet states, **1d**, **2d**, and **3d**, respectively. The lowest C_2Tl_4 triplet state lies between **4c** and **4d** and is 0.1 kcal/mol higher in energy than **4c**. Therefore, we may safely conclude that the singlet species, **1a–4a** shown in Figure 1, are indeed the global minima under the considerations of both singlet and triplet states.

We noted that a time-of-flight mass spectrum from laser vaporization of a mixed graphite/aluminum target, reported by Boldyrev–Wang groups,⁶³ showed a strong peak on the $C_2Al_4^-$ component, which may relate to the fact that C_2Al_4 is the global minimum. We propose further verifications of the species and others reported.

The B3LYP and MP2 optimizations gave the same structures for these isomers except for **1b** (see Figure 2). The first three isomers in the C_2Ga_4 (**2b–2d**) and C_2In_4 (**3b–3d**) series have similar structures and the same energy order, but **1d** and **4c** have no counterparts in the C_2Ga_4 and C_2In_4 series. It is remarkable because they are ranked entirely based on the random searches, which implies some common features in their chemical bonding. For all of the four series, the relative energy orders of the isomers, given by MP2 and CCSD(T)//MP2, are identical in spite of their different magnitudes. For C_2Al_4 series, the B3LYP order is consistent with those given by MP2 and CCSD(T)//MP2, but the two levels give different orders for C_2Ga_4 , C_2In_4 , and C_2Tl_4 series and reverse the orders of **c** and **d** isomers. Nevertheless, the three levels unanimously predicted **a** and **b** isomers to be the first two lowest isomers.

The C_{2h} C_2Be_4 (**M2**) with dptCs is a local minimum. The exploration on its PES, at the B3LYP/6–31+G* level, indicates that it is the fifth lowest. The refinement calculations at the B3LYP/aug-cc-pVTZ also rank it the fifth lowest. But the calculations at both MP2 and CCSD(T)//MP2 levels rank it the sixth lowest. At the three levels, the dptC-contained C_2Be_4 is 8.8, 11.9, and 9.7 kcal/mol higher in energy than the global minima (**5a**), respectively (see Figure 3, their Cartesian coordinates are given in SI2). It should be mentioned that PES exploration (B3LYP/6–31+G*) revealed that the global minimum is a planar irregular hexagon. In a recent B3PW91/6–31+G* study on the Be_nC_m ($n = 1\sim 10; m = 1\sim 11 - n$, respectively) clusters, Mainardi and co-workers⁶⁴ also reported such a hexagonal structure. However, the B3LYP/aug-cc-pVTZ optimization indicates that it is the regular hexagonal global minimum with C_{2v} symmetry. The MP2 optimizations with aug-cc-pVTZ basis set also led to the C_{2v} structure but rank it the third lowest. However, the CCSD(T)//MP2 single-point calculations again indicate the C_{2v} structure to be the global minimum.

Among the twelve vibrational modes of **1a–4a**, only the asymmetric bridging-E(E_{BR})–C and terminal-E(E_{TE})–C stretching modes show strong activities and the others are entirely inactive or have very small infrared intensities. To aid the experimental identification, Table 1 lists the frequencies and infrared intensities of the two strongest modes of **1a–4a**, along with the strongest modes of CaI_4^{2-} and CaI_5^+ . Although the irreducible representations of C–E stretching vibrational modes at B3LYP and MP2 level for C_2Ga_4 , C_2In_4 , and C_2Tl_4 are

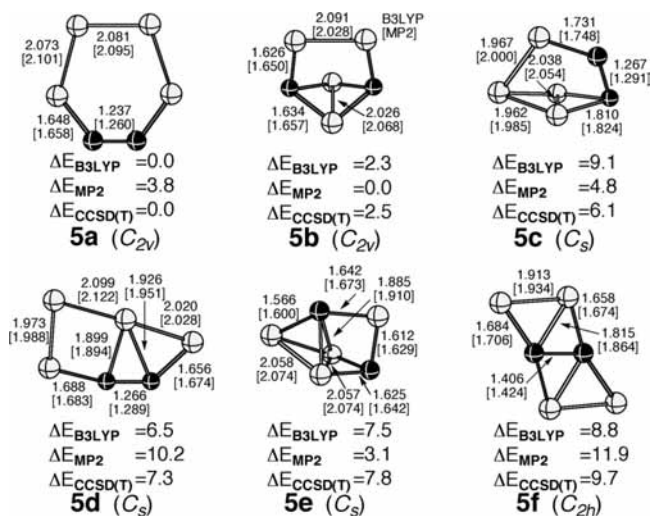


Figure 3. B3LYP and MP2 optimized geometries of the six lowest isomers of C_2Be_4 , together with the relative energies (ΔE in kcal/mol) and the key bond lengths (in Å). Values in brackets are MP2 results.

different due to their different structures, we have confirmed that the motions of the modes are the same. This holds for C_2-E_{BR} vibrational modes of the C_2Ga_4 vs C_2E_4 ($E = In$ and Tl) at the MP2 level. Similar to **1a–4a**, the infrared intensities of other modes of CAI_4^{2-} and CAI_5^+ are much smaller. Therefore, their infrared spectra only have strong peaks on the two modes, which may facilitate experimental identification of these species. As shown in the Table 1, the B3LYP frequency values of **1a–4a** ($326.8\sim 547.3\text{ cm}^{-1}$ for $E_{BR}-C$ stretching, $385.0\sim 602.4\text{ cm}^{-1}$ for $E_{TE}-C$ stretching) are in agreement with those of MP2 values. The decreasing vibrational frequencies from **1a** to **4a** indicate the weakening C–E bonding as E moves down the periodical table. Note that the vibration frequencies of $Al_{TE}-C$ and $Al_{BR}-C$ stretching in **1a**, 602.4 and 547.3 cm^{-1} , respectively, are comparable to the Al–C stretching at 736.1 cm^{-1} for CAI_4^{2-} and 631.9 cm^{-1} for CAI_5^+ , respectively. Consistently, the $Al_{TE}-C$ and $Al_{BR}-C$ bond lengths, 1.994 and 2.155 Å , respectively are compared with the 1.977 Å in CAI_4^{2-} and 2.107 Å in CAI_5^+ .

CAI_4^{2-} and CAI_5^+ are isoelectronic and have been characterized to be the global minima.^{27,32} In their common C_{2v} subgroup, CAI_4^{2-} and CAI_5^+ have the same electron configurations except for the orders of the three highest occupied molecular orbitals (MOs) (eqs 2 and 3). C_2Al_4 has two electrons more than CAI_4^{2-} and CAI_5^+ and the two extra electrons occupy HOMO-2 (framed in eq 1) which is absent in CAI_4^{2-} and CAI_5^+ . Except for this orbital, the occupied MOs of **1a** have the same symmetries (orbital shapes) as those of CAI_4^{2-} and CAI_5^+ (eqs 1–3 and Figure 4). The HOMO of CAI_4^{2-} and the HOMO-2 of CAI_5^+ ,

which were proposed to contribute to stabilizing their ptC arrangements, have a counterpart in **1a** (i.e., the HOMO). Note that a similar orbital is also present in the boraplanes designed by Wang and Schleyer.⁶⁵ The Wiberg bond indices (WBIs) and the NBO charges given by NBO analyses are summarized in Table 2. Although **1a**, CAI_4^{2-} , and CAI_5^+ are in the different charge states, the C_2 moiety in **1a** bears charges ($-2.78e$), close to those of the carbon centers in CAI_4^{2-} ($-2.72e$) and CAI_5^+ ($-2.80e$), which, along with their similarities in MO shapes (see Figure 4), suggests that the C_2 moiety in **1a** can be viewed as the equivalences of the carbon centers in CAI_4^{2-} and CAI_5^+ . If the C_2 moiety in **1a** are viewed as an unit to interact with the peripheral Al_4 ring, then its total WBI of the moiety, $2.02 [(3.11-2.10) \times 2]$, to the Al_4 periphery is comparable to those, 2.24 in CAI_4^{2-} and 2.15 in CAI_5^+ . This further rationalizes our deduction. Recall that the C–Al stretching vibrational modes in **1a**, CAI_4^{2-} , and CAI_5^+ all are most strongly active while other modes are inactive or have negligible infrared intensities. Therefore, the bonding interactions between the C_2 moiety and the Al_4 periphery in **1a** are similar to those between C and Al_4 or Al_5 periphery in CAI_4^{2-} and CAI_5^+ , respectively. Indeed, the WBIs of $Al_{BR}-$ and $Al_{TE}-C_2$ moiety in **1a**, 0.51 and 0.56 , respectively, are comparable to the 0.56 in CAI_4^{2-} and 0.43 in CAI_5^+ (Table 2). Due to their different charge states, the electrostatic interactions between the centers and the Al_4 (or Al_5) periphery vary, as reflected by the NBO charges on the Al atoms. However, the Al–Al bonding in **1a** are different from those in CAI_4^{2-} and CAI_5^+ . The sum of WBIs of $Al_{BR}-Al$ (0.29) and $Al_{TE}-Al$ (0.22) are significantly smaller than the 1.34 in CAI_4^{2-} and 1.39 in CAI_5^+ . Among the occupied MOs, one can identify that the major contributions to the differences come from HOMO-6 of **1a**. This MO in **1a** is involved in C– Al_{TE} and C–C σ bonding, whereas the corresponding orbital (HOMO-5) in CAI_4^{2-} and CAI_5^+ contribute to the σ Al peripheral bonding. Because **1a** is neutral, the charges on the C_2 moiety are donated by the peripheral Al atoms, the Al_{BR} and Al_{TE} bear $0.69e$ and $0.70e$ positive charges, respectively. In contrast, the Al atoms in CAI_4^{2-} have a much smaller positive charge ($0.18e$), which is due to the net negative $2e$ charges of the dianion. Although CAI_5^+ possesses $1e$ positive charge, the Al atoms in CAI_5^+ have $0.76e$ positive charges, comparable to those in **1a**. This is because the extra Al atom in CAI_5^+ compensates the effects of the monocationic charge state of CAI_5^+ .

$$C_2Al_4 : [Core]A_1^2B_2^2A_1^2A_1^2B_1^2A_1^2B_2^2 \boxed{A_1^2} B_2^2A_1^2 \quad (1)$$

$$CAI_4^{2-} : [Core]A_1^2B_2^2A_1^2A_1^2B_1^2A_1^2B_2^2 \quad B_2^2A_1^2 \quad (2)$$

$$CAI_5^+ : [Core]A_1^2B_2^2A_1^2A_1^2B_1^2A_1^2 \quad B_2^2B_2^2 \quad (3)$$

The natural electron configurations of Al atoms in **1a**, CAI_4^{2-} , and CAI_5^+ are compared in Table 3. Similar to CAI_4^{2-} and

TABLE 1: Wavenumbers (WN, in cm^{-1}) of the Most Active C–E Stretching Modes with Their Symmetries (SYM) and Intensities (Int, in km/mol) of **1a–4a, CAI_4^{2-} and CAI_5^+ , Calculated at B3LYP and MP2 (in parentheses) Levels^a**

	$E_{TE}-C$ stretching			$E_{BR}-C$ stretching		
	SYM	WN	Int	SYM	WN	Int
C_2Al_4 (1a)	B_{1u} (B_{1u})	602.4 (615.9)	635.2 (626.7)	B_{2u} (B_{2u})	547.3 (571.6)	372.3 (366.9)
C_2Ga_4 (2a)	B_{1u} (B_2)	499.0 (581.2)	527.8 (617.0)	B_{2u} (A_1)	423.7 (539.0)	252.3 (319.2)
C_2In_4 (3a)	B_{1u} (B_2)	422.4 (497.2)	352.9 (450.7)	B_{2u} (B_1)	370.6 (458.1)	172.9 (256.3)
C_2Tl_4 (4a)	B_{1u} (B_2)	385.0 (468.8)	229.4 (397.5)	B_{2u} (B_1)	326.8 (439.3)	93.4 (210.2)
CAI_4^{2-}	E_u (E_u)	736.1 (785.2)	652.1 (866.6)			
CAI_5^+	E_1' (E_1')	631.9 (626.5)	380.0 (344.9)			

^a Note that the C–Al stretching modes of CAI_4^{2-} and CAI_5^+ are degenerated.

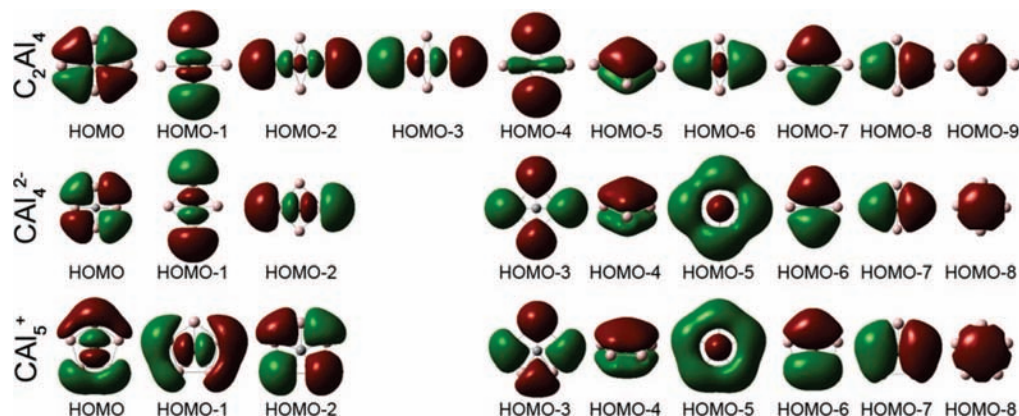


Figure 4. Comparisons of the occupied valence MOs of **1a** with those of CAI_4^{2-} and CAI_5^+ .

TABLE 2: NBO Charges (in e) and Wiberg Bond Indices (WBI) of **1a–4a**, CAI_4^{2-} , and CAI_5^+

	q_C	q_E		WBI _C	WBI _E		WBI _{C–C}	WBI _{E–CC}		WBI _{E–Es}	
		E_{BR}	E_{TE}		E_{BR}	E_{TE}		E_{BR}	E_{TE}	E_{BR}	E_{TE}
C_2Al_4 (1a)	–1.39	+0.69	+0.70	3.11	0.80	0.72	2.10	0.51	0.50	0.29	0.22
C_2Ga_4 (2a)	–1.19	+0.59	+0.60	3.30	0.86	0.83	2.17	0.57	0.54	0.32	0.25
C_2In_4 (3a)	–1.11	+0.55	+0.56	3.36	0.90	0.89	2.35	0.48	0.54	0.42	0.35
C_2Tl_4 (4a)	–1.01	+0.50	+0.51	3.42	0.93	0.95	2.44	0.45	0.53	0.48	0.42
C_2B_4	–1.02	+0.45	+0.57	3.43	1.31	0.90	1.74	0.95	0.73	0.36	0.16
CAI_4^{2-}	–2.72		+0.18	2.24		1.90			0.56		1.34
CAI_5^+	–2.80		+0.76	2.14		1.99			0.43		1.56
C_2H_4	–0.37		+0.18	3.94		0.81	2.05		0.95		0.02

TABLE 3: Natural Electron Configurations of E Atoms in **1a–4a**, CAI_4^{2-} , and CAI_5^+

		s	p_x	p_y	p_z
1a	Al_{TE}	1.83	0.26	0.19	0.01
	Al_{BR}	1.83	0.24	0.22	0.01
2a	Ga_{TE}	1.88	0.27	0.24	0.01
	Ga_{BR}	1.90	0.30	0.19	0.02
3a	In_{TE}	1.90	0.23	0.34	0.01
	In_{BR}	1.92	0.35	0.35	0.01
4a	Tl_{TE}	1.93	0.22	0.34	0.01
	Tl_{BR}	1.95	0.39	0.15	0.01
CAI_4^{2-}		1.53	0.61	0.56	0.10
CAI_5^+		1.45	0.38	0.30	0.09

CAI_5^+ , the Al $3p_z$ orbital in **1a** has very small occupancy, $0.01e$, in comparison with $0.10e$ (CAI_4^{2-}) and $0.09e$ (CAI_5^+). Because CAI_4^{2-} is $2e$ negatively charged, its σ electron ($3p_x + 3p_y$) occupancy ($1.17e$) of Al atom is larger than the $0.68e$ in CAI_5^+ and $0.45e$ of $\text{Al}_{\text{TE}}/0.46e$ of Al_{BR} . The natural electron configuration points out that Al ligands mainly serve as σ electron donors in the three species. Previously, on the basis of the electronic structures of planar methane, Hoffmann, Alder, and Wilcox¹ suggested using σ -donor/ π -acceptor ligands to enhance the electronic deficient σ bonding between ptC and ligands and delocalize the unfavorable lone pair on ptC of planar methane. However, CAI_4^{2-} and CAI_5^+ have MOs similar to the lone pair in the planar methane. Note that there are π interactions between ptC and peripheral Al atoms, but the interactions are small (see HOMO-4 in Figure 4). We⁶⁶ have rationalized the phenomena in the $\text{C}(\text{BeH})_4^{2-}$ system, and partially attributed to the weakening effects on the rigidity of the tetrahedral carbon, exerted by the metallic ligands. In contrast, **1a** has no such π lone-pair but a C–C π bonding orbital (HOMO-5) predominantly located on the C_2 moiety. One may ask that the Al–Al bonding in CAI_4^{2-} and CAI_5^+ is much stronger than that in **1a**, why **1a** still can achieve the planarity? The answer is that the energetic benefit due to the C–C π bonding in **1a** compensate the energetic favorable Al_4 (or Al_5) σ peripheral bonding

(HOMO-5) in CAI_4^{2-} and CAI_5^+ , which are supported by the relatively larger WBI_C (3.11 vs. 2.24 and 2.14) and smaller WBI_{Al–Als} (less than 0.3 vs. 1.34 and 1.56) in **1a** than those in CAI_4^{2-} and CAI_5^+ .

The above analyses have revealed that the C_2 moiety in **1a** can be considered as the equivalence of carbon center in CAI_4^{2-} and CAI_5^+ . How feasible can another carbon be embedded into the Al_4 ring? If considering **1a** as the product of CAI_4 (T_d , the global minimum⁶⁷) + C (3P) reaction, the reaction enthalpies (ΔH°) are -156.9 , -163.8 , -185.0 kcal/mol at the B3LYP, MP2, and CCSD(T)/MP2 levels, respectively, indicating the high energetic favorableness of incorporating a C atom into T_d CAI_4 .

Summarizing the above comparisons, one can concisely understand the stabilities of these species as illustrated by Figure 5. For CAI_4^{2-} , addition of the two electrons in T_d CAI_4 results in the ptC global minimum, because the two extra electrons take the advantage of the peripheral bonding MO²⁷ (a useful orbital not being utilized in neutral D_{4h} CAI_4 , which is the third-order saddle point). The two electrons in CAI_5^+ are introduced by adding an Al^+ in the Al_4 ring. In **1a**, the addition of two electrons is fulfilled by introducing another carbon in the center. Although the carbon brings four electrons, two of them are used for an internal C–C bonding. The neutrality of **1a** may facilitate

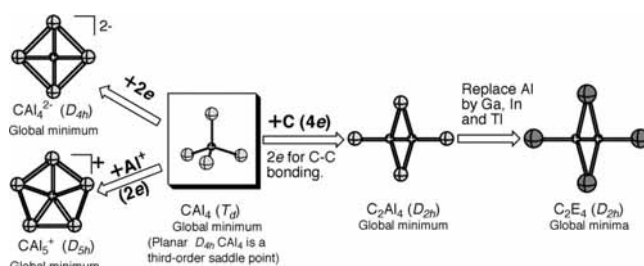


Figure 5. Relationships of CAI_4 , CAI_4^{2-} , CAI_5^+ , and C_2Al_4 .

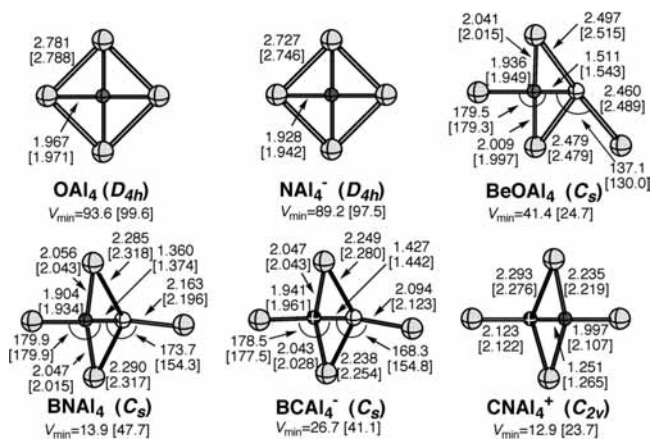


Figure 6. B3LYP and MP2 optimized geometries of OAl₄, NAl₄⁻, BeOAl₄, BAl₄, CAl₄⁺, and BCAl₄⁻, together with the smallest harmonic frequencies (V_{\min} in cm^{-1}) and the geometrical parameters (the bond lengths in Å and bond angles in degree. The values in bracket are the MP2 results).

generation and detection in the experiments (e.g., low temperature matrix isolation).

The established equivalence relationship can be used to design more molecules with planar tetracoordinate heteroatoms (Figure 6, the Cartesian coordinates are given in the SI3). The previously reported OAl₄ and NAl₄⁻ are reconfirmed to be minima in the present levels of theories. Similarly, at both B3LYP and MP2 levels, the replacements of C₂ moiety in C₂Al₄ by the isoelectronic BeO, BN, BC⁻, and CN⁺ moieties also lead to minima, BeOAl₄, BAl₄, BCAl₄⁻, and CAl₄⁺. Due to the stronger metallic characters and bigger radius of Be and B than C, the terminal Al atoms bonded to Be or B in BeOAl₄ and BAl₄ and BCAl₄⁻ bend away from the BeO or BC axes. Consistently, the bending Al atom in BeOAl₄ is further away from the BeO axis than the Al atoms in BAl₄ and BCAl₄⁻ from BN and BC axis. The terminal Al atoms in CAl₄⁺ lie along the CN axis.

The C–C double bonding ($\text{WBI}_{\text{C}=\text{C}} = 2.11$) and *D_{2h}* symmetry of **1a** reminds us of the classical ethene. As reflected by the WBIs and NBO charges (Table 2), the replacement of hydrogens in ethene by alumina results in the common differences expected on the basis of the fact that Al is more metallic than H. The MO analyses give more insights. As compared in the Figure 7, the eight more valence electrons in **1a** than in C₂H₄ can be assigned to the HOMO-4 to HOMO-1 MOs (right column), which can be considered as nonbonding MOs. Note that CAl₄²⁻ and CAl₅⁺ also have four occupied nonbonding MOs (including the HOMO-5, see Figure 4). Interestingly, although Al has two more valence electrons available than hydrogen, they are not efficiently utilized for bonding interactions and occupy the nonbonding orbitals (i.e., the HOMO-1~HOMO-4 of C₂Al₄, as shown in Figure 7) in these species. The pair of HOMO-9 and HOMO-8 in **1a** corresponds to the pair of HOMO-5 and HOMO-4 in C₂H₄, which originates from the bonding and antibonding interactions between two carbon 2s orbitals. The π (HOMO-5) and σ (HOMO-6) MOs in **1a** have counterparts in ethene, HOMO and HOMO-2, respectively. The essential differences lie in the HOMO and HOMO-7 of **1a**. On the basis of the consideration of orbital symmetry, the two MOs of **1a** correspond to HOMO-1 and HOMO-3 of ethene, respectively. The ethene HOMO-1 contributes to the C–H σ bonding and has negative contribution to the C–C bonding, whereas the corresponding orbital (HOMO) of **1a** contributes to the peripheral Al₄ bonding and the antibonding interaction between carbons is negligible, which reflects the differences between

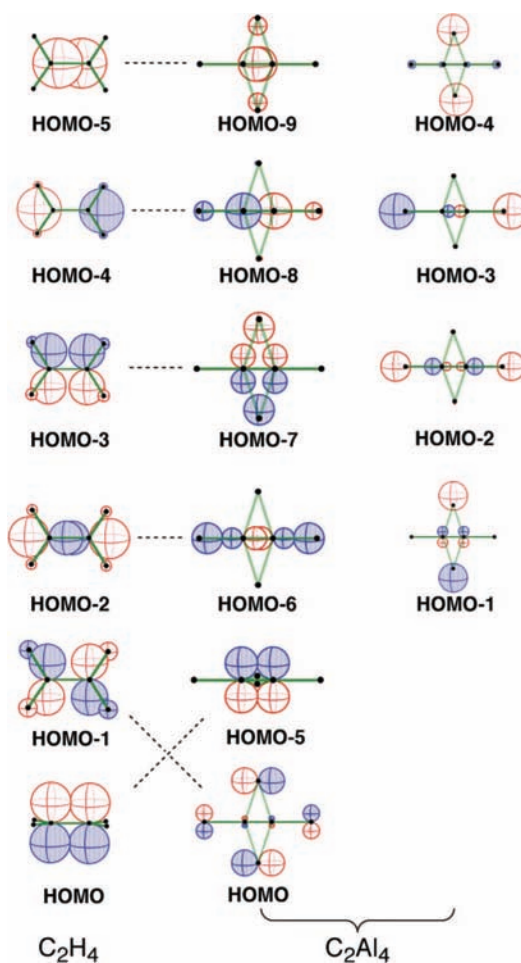


Figure 7. Comparison of the valence MOs of **1a** with those of ethylene.

Al and H atoms, that is, Al has p electrons available for such covalent bonding. The ethene HOMO-3 contribute to the four C–H bonds and C–C bonding, in contrast, the electrons in HOMO-7 of **1a** bound two of the Al atoms in the bridging positions. This can be attributed to the electron deficiency of the Al atom, similar to the bridging hydrogens in carbocations.

The MOs of **1a**, **2a**, **3a**, and **4a**, given in SI6, indicate the similarities in their electronic structures, although the MO energy order may vary. However, the individual characters of the ligands results in differences in their bonding as quantified by WBI and NBO charges. As the E atoms move down the periodical table (from Al to Tl), because their electronegativities (1.5(Al), 1.6(Ga), 1.7(In), 1.8(Tl))⁶⁸ increase, the NBO charges on the C₂ moieties, -2.78 , -2.38 , -2.22 , and $-2.02e$ for **1a–4a**, respectively, decrease gradually. The decreasing C–C bond lengths, 1.311, 1.298, 1.282, and 1.274 Å in **1a–4a**, respectively, are consistent with the increasing $\text{WBI}_{\text{C}=\text{C}}$, 2.10(**1a**), 2.17 (**2a**), 2.35 (**3a**), and 2.44 (**4a**), respectively.

Boron is the only exception for Group IIIA elements to form *D_{2h}* C₂B₄ species with dptCs. Zeng et al.³¹ have explored the PES of C₂B₄. Our calculations at the B3LYP/aug-cc-pVTZ level showed the *D_{2h}* C₂B₄ with dptCs is a forth-order saddle point and 187.3 kcal/mol higher in energy than the reported global minimum,³¹ in which two carbons bond to a B₄ ring externally. We reason as follows: As discussed in our C(BeH₄)₄²⁻ paper,⁶⁶ the rigidity of a tetrahedral carbon become weaker as the ptC-ligand bonds become less covalent. The average WBI of C–B bonds in the *D_{2h}* C₂B₄, 0.84, is substantially larger than the average WBIs (0.51, 0.56, 0.51, and 0.49) in **1a–4a**, respec-

tively. Although D_{2h} C_2B_4 also possesses a B_4 peripheral occupied MOs similar to HOMO and HOMO-5 of C_2Al_4 , the B_4 ring is too small to accommodate two carbons. Note that the D_{4h} CB_4^{2-} is even not a minimum (it is a second-order saddle point).

3. Conclusions

Using ab initio and DFT calculations, we have searched for the smallest species with dptCs. The scanning on the candidates, C_2E_4 (E = the second- and third-row main group elements) resulted in three minima (C_2E_4 (E = Be, Mg, and Al)). The DFT PES exploration shows the C_2E_4 (E = Al–Ti) not only has structures with dptCs but are also global minima. At the CCSD(T)//MP2 level, they are confirmed again to be the global minima, and locate 18.0, 18.3, 13.4, and 12.2 kcal/mol lower than their nearest isomers, respectively. C_2Be_4 is a local minimum. The comparisons of D_{2h} C_2Al_4 with CAI_4^{2-} and CAI_5^+ uncover that the C_2 moiety in the former is the equivalence of the carbon centers in the later. Concisely, the stabilities of these species can be rationalized as follow: The addition of two electrons in CAI_4 (i.e., CAI_4^{2-}) can utilize a Al_4 peripheral bonding MO. CAI_5^+ obtains the two electrons by incorporating an Al^+ atom in the ring. The addition of two electrons in C_2Al_4 (and its analogs) is fulfilled by incorporating a carbon in the center of CAI_4 . Although the carbon brings four electrons, two of them are utilized for an internal C–C bonding. Although DFT calculations characterized **1a–4a** to have D_{2h} structures with perfect dptCs, the MP2 optimizations distorted the D_{2h} symmetry slightly. The energy differences between D_{2h} and the distorted structures are marginal (<0.02 kcal/mol). The MO comparisons of **1a** with ethene further rationalized the origins of **1a–4a** to achieve dptCs.

Acknowledgment. We dedicate the work to Professor Pin Yang on the occasion of his 75th birthday. The authors thank the Shanxi University, the Chinese Academy of Sciences and NSFC (20773160 to Z.X.W. and 20573088 to S.D.L.) for funding. The authors appreciate Prof. Yue-Kui Wang and Dr. Cai-Xia Yuan for their valuable discussions.

Supporting Information Available: The Cartesian coordinates and total energies of **1a–4a** and their three singlet isomers, three triplet isomers, the C_2Be_4 isomers and the OAl_4 , NAI_4^- , $BeOAl_4$, $BNAI_4$, $BCAl_4^-$, and $CNAI_4^+$, the pictures and relative energies of triplet isomers for **1a–4a**, the pictures of occupied valence MOs of **1a–4a** and C_2B_4 , and the complete ref 60. This material is available free of charge via the Internet at <http://pubs.acs.org>.

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JP8099187